A STUDY ON MHD BOUNDARY LAYER FLOW OVER A NONLINEAR STRETCHING SHEET USING IMPLICIT FINITE DIFFERENCE METHOD

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Abstract

The problem of the MHD boundary layer flow of an incompressible viscous fluid over a non-linear stretching sheet is discussed. Using similarity transformation governing equations are transformed into nonlinear ordinary differential equation. The Implicit finite difference Keller box method is applied to find the numerical solution of nonlinear differential equation. Graphical results of fluid velocity have been presented and discussed for the related parameters.

Keywords: MHD, boundary layer, nonlinear differential equations, and Keller box method

1. INTRODUCTION

In 1973, McCormack & Crane [8] introduced the stretching sheet problem. The stretching problems for steady flow have been used in various engineering and industrial processes, like non-Newtonian fluids, MHD flows, porous plate, porous medium and heat transfer analysis.

Magnetohydrodynamics (MHD) is the study of the interaction of conducting fluids with electromagnetic phenomena. The flow of an electrically conducting fluid in the presence of a magnetic field is important in various areas of technology and engineering such as MHD power generation, MHD flow meters, MHD pumps, etc.


Rashidi in [4] used the differential transform method with the Padé approximant and developed analytical solutions for this problem. In [1, 2] applied the HAM in order to obtain an analytical solution of the governing nonlinear differential equations.

In this paper, MHD boundary layer equation is solved with the help of implicit finite difference Keller box method and various results are discussed graphically.

2. GOVERNING EQUATIONS

Consider the Magnetohydrodynamic flow of an incompressible viscous fluid over a stretching sheet at $y = 0$. The fluid is electrically conducting under the influence of an applied magnetic field $B(x)$ normal to the stretching sheet. The induced magnetic field is neglected. The resulting boundary layer equations are as follows [6]

\begin{equation}
\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0
\end{equation}

\begin{equation}
u \frac{\partial u}{\partial x} + \nu \frac{\partial v}{\partial y} = \nu \frac{\partial^2 u}{\partial y^2} - \frac{\sigma B^2(x)}{\rho} u
\end{equation}

where $u$ and $v$ are the velocity components in the $x$ - and $y$ - directions respectively, $\nu$ is the kinematic viscosity, $\rho$ is the fluid density and $\sigma$ is the electrical conductivity of the fluid. In equation (2), the external electric field and the polarization effects are negligible and following Chiam [7] we assume that the magnetic field $B$ takes the form

\[ B(x) = B_0 x^{(n-1)/2} \]

The boundary conditions corresponding to the non-linear stretching of the sheet are [6]

\[ u(x,0) = cx^n, v(x,0) = 0 \]
\( u(x, y) \to 0, y \to \infty \) where c and n are constants.

Using the following substitutions

\[
  t = \sqrt[2n]{\frac{c(n+1)}{2y}} x^{\frac{n}{2}} y
\]

\[
  u = cx^n f(t)
\]

\[
  v = -\sqrt[2n]{\frac{c(n+1)}{2y}} x^{\frac{n}{2}} \left[ f(t) + \frac{n-1}{n+1} t f'(t) \right]
\]

Equation (1)-(2) are transformed to

\[
  f''(t) + f(t) f'(t) - \beta \left( f'(t) \right)^2 - Mf'(t) = 0 \tag{3}
\]

With the following boundary conditions

\[
  f(0) = 0, f'(0) = 1, f'(\infty) = 0 \tag{4}
\]

\[
  \beta = \frac{2n}{n+1}, M = \frac{2\sigma B^2}{\rho c(1+n)}
\]

3. KELLER BOX METHOD

Equation (3) subject to the boundary conditions (4) is solved numerically using implicit finite difference method that is known as Keller-box in combination with the Newton’s linearization techniques as described by Cebeci and Bradshaw [3]. This method is unconditionally stable and has second-order accuracy.

In this method the transformed differential equation (3) are written in terms of first order system, for that introduce new dependent variable \( u, v \) such that

\[
  f' = u \tag{5}
\]

\[
  u' = v \tag{6}
\]

Where prime denotes differentiation with t.

\[
  v + fv - \beta u^2 - Mu = 0 \tag{7}
\]

With the new boundary conditions

\[
  f(0) = 0, u(0) = 1, u(\infty) = 0 \tag{8}
\]

Now, write the finite difference approximations of the ordinary differential equations (5)-(6) for the midpoint \( x^n, j_{-\frac{1}{2}} \) of the segment using centered difference derivatives, this is called centering about \( x^n, j_{-\frac{1}{2}} \).

\[
  f^n_j - f^n_{j-1} \left/ h_j \right. = u^n_j + u^n_{j-1} \left/ 2 \right. \tag{9}
\]

\[
  u^n_j - u^n_{j-1} \left/ h_j \right. = v^n_j + v^n_{j-1} \left/ 2 \right. \tag{10}
\]

Now, ordinary differential equation (7) is approximated by the centering about the mid-point \( x^n, j_{-\frac{1}{2}} \) of the rectangle.

In first step we center about \( j \) for simplicity, remove n.

\[
  \left( v^n_j - v^n_{j-1} \right) + \left[ f_j^n + f_{j-1} \right] \left/ 2 \right. \left[ v^n_j + v^n_{j-1} \right] = \beta \left[ u^n_j + u^n_{j-1} \right] - M \left[ u^n_j + u^n_{j-1} \right] =
\]

\[
  \left[ u^n_j + u^n_{j-1} \right] \left/ 2 \right. \left[ v^n_{j+1} - v^n_{j-1} \right] = \beta \left( u^n_j + u^n_{j-1} \right) - M \left( u^n_j + u^n_{j-1} \right) =
\]

\[
  \left[ u^n_j + u^n_{j-1} \right] \left/ 2 \right. \left[ v^n_{j+1} - v^n_{j-1} \right] = \beta \left( u^n_j + u^n_{j-1} \right) - M \left( u^n_j + u^n_{j-1} \right) =
\]

Now linearize the nonlinear system of equations (9)-(11) using the Newton’s quasi-linearization method [5].

For that use,
\[ f^{n+1}_j = f^n_j + \delta f^n_j \]
\[ u^{n+1}_j = u^n_j + \delta u^n_j \]
\[ v^{n+1}_j = v^n_j + \delta v^n_j \]

Equation (8) to (10) can be rewritten as

\[ (\delta f_j - \delta f_{j-1} - \frac{h_j}{2}(\delta u_j + \delta u_{j-1}) = (r^j_1) \] (12)
\[ (\delta u_j - \delta u_{j-1} - \frac{h_j}{2}(\delta v_j + \delta v_{j-1}) = (r^j_2) \] (13)
\[ (a_1)\delta v_j + (a_2)\delta v_{j-1} + (a_3)\delta f_j + (a_4)\delta f_{j-1} + (a_5)\delta u_j + (a_6)\delta u_{j-1} = (r^j_3) \] (14)

The linearized difference system of equations (12) - (14) has a block tridiagonal structure. In a vector matrix form, it can be written as

\[
\begin{bmatrix}
[A_1] & [C_1] \\
[B_1] & [A_1] & [C_1] \\
\vdots & \vdots & \vdots \\
[B_{j-1}] & [A_{j-1}] & [C_{j-1}] \\
[B_j] & [A_j] & [C_j] \\
\end{bmatrix}
\begin{bmatrix}
[\delta f_j] \\
[\delta v_j] \\
\vdots \\
[\delta f_{j-1}] \\
[\delta v_{j-1}] \\
\end{bmatrix} =
\begin{bmatrix}
[r^j_1] \\
[r^j_2] \\
\vdots \\
[r^j_3] \\
[r^j_3] \\
\end{bmatrix}
\]

This block tridiagonal structure can be solved using LU method explained by Na [5].

4. RESULT AND DISCUSSION

Graphically, effects of magnetic parameter and non-dimensional parameter \( \beta \) on velocity profile are shown in following figures. Fig.1 and 2 shows, as the magnetic parameter M and non-dimensional parameter \( \beta \) increases, the fluid velocity decreases. Also it can be observed that as magnetic parameter M increases velocity profile decreases very rapidly as compared to increase of non-dimensional parameter \( \beta \). Fig 3 and 4 shows the effect of magnetic parameter on \( f \).

**Fig -1:** Effect of Magnetic Parameter on fluid velocity obtained by Keller Box Method for \( \beta = 1 \)
Fig -2: Effect of $\beta$ on fluid velocity obtained by Keller Box Method when $M = 2$

Fig -3: Effect of Magnetic Parameter on $f(t)$ obtained by Keller Box Method for $\beta = 1$

Fig -4: Effect of Magnetic Parameter on $f(t)$ obtained by Keller Box Method for $\beta = -1$
CONCLUSIONS

In this study, MHD viscous flow over a stretching sheet is considered. Keller Box Method is applied to solve the governing nonlinear differential equation. The results indicate that the velocity will be reduced by increasing the two parameters of M and $\beta$; however, effect of “M” in reduction of the velocity components is more than “$\beta$”. Therefore, it can be concluded that Keller Box is one of the best method to study on MHD viscous flow numerically and get the appropriate results.

REFERENCES


